

Recall:

• IVP
$$\begin{cases} \dot{u}(t) + a u(t) = f(t) \\ u(0) = u_0 \end{cases}$$

• VOC + Stability for $u(t)$

• c(1) $\leadsto u(t) = u_0 \frac{t_1 - t}{t_1} + u(t_1) \frac{t}{t_1}$ for $t \in [0, t_1]$

where
$$\left(1 + \frac{a t_1}{2}\right) u(t_1) = \left(1 - \frac{a t_1}{2}\right) u_0 + \int_0^{t_1} f(\tau) d\tau$$

+ appropriate quadrature formula gives C-N

• d(0) $\leadsto u(t) = u_n = u_{n-1} - \frac{1}{2} a u_n + \int_{t_{n-1}}^{t_n} f(\tau) d\tau$ for $t \in (t_{n-1}, t_n]$

+ appropriate quadrature formula gives back. Euler

5) A priori error estimates for dG(1):

Prob:
$$\begin{cases} \dot{u}(t) + a u(t) = f(t) & 0 < t \leq T \\ u(0) = u_0 \end{cases}$$

Then Denote by $U(t)$ the approximation of $u(t)$ given by dG(1), denote by k_n the length of the small interval $[t_{n-1}, t_n]$.

If $|k_n| \leq 1/2$ for $n=1, 2, 3, \dots$, then

$$|u(T) - U(T)| \leq \frac{e}{4} (e^{2|a|T} - 1) \max_{1 \leq n \leq N} (k_n \max_{t \in [t_{n-1}, t_n]} |\ddot{u}(t)|)$$

Error $\sim O(k^2)$ when one has uniform $k_n = k$.

Chapter VIII: The heat equation in 1d

Goal: Discuss stability of exact sol.

Present a numerical method for heat eq.

1) Heat equations in 1d:

The following is a simple model to describe

heat flow on a thin wire/metal rod of length 1:

(Newton's law cooling)

$$(M) \begin{cases} u_t(x,t) - u_{xx}(x,t) = f(x,t) & 0 < x < 1, 0 < t \leq T \quad (DE) \\ u(0,t) = 0, u_x(1,t) = 0 & 0 < t \leq T \quad (BC) \\ u(x,0) = u_0(x) & 0 < x < 1 \quad (IC) \end{cases}$$

Here; $u(x,t)$... unknown $\leadsto$ temperature at position x and time t

$f(x,t)$... given $\leadsto$ heat source/sink

$u_0(x)$... given $\leadsto$ initial temperature profile

$u(0,t) = 0$... homog. Dirichlet BC
 $\leadsto$ fixed temp. 0 at $x=0$

$u_x(1,t) = 0$... homog. Neumann BC
 $\leadsto$ no heat flux at $x=1$.

Th; The sol. to (M) satisfies the following stability estimates:

$$(i) \|u(\cdot, t)\|_{L^2(0,1)} \leq \|u_0\|_{L^2(0,1)} + \int_0^t \|f(\cdot, s)\|_{L^2(0,1)} ds$$

(remember last chapter)

$$(ii) \|u_x(\cdot, t)\|_{L^2(0,1)} \leq \|u_0'\|_{L^2(0,1)} + \int_0^t \|f(\cdot, s)\|_{L^2(0,1)} ds$$

(iii) When $t \geq 0$, one has

$$\|u(\cdot, t)\|_{L^2(0,1)} \leq \|u_0\|_{L^2(0,1)} \cdot e^{-2t} \quad (\text{typical for parabolic})$$

Rem! • Notation: $\|u(\cdot, t)\|_{L^2(0,1)}^2 = \int_0^1 |u(x, t)|^2 dx$

• (iii) says that the temperature goes to zero in L^2 -norm as $t \rightarrow \infty$.

Proof,

(i) Multiply (DE) with u and integrate $\int_0^1 dx$:

$$\underbrace{\int_0^1 u_t(x, t) u(x, t) dx}_{\text{by part}} - \underbrace{\int_0^1 u_{xx}(x, t) u(x, t) dx}_{\text{by part}} = \int_0^1 f(x, t) u(x, t) dx$$

$$\frac{1}{2} \frac{\partial}{\partial t} (u(x, t))^2 = \frac{2}{2} u(x, t) \frac{\partial}{\partial t} u(x, t)$$

(chain rule)

Integrate by part:

$$\frac{1}{2} \frac{d}{dt} \underbrace{\int_0^1 (u(x, t))^2 dx}_{\|u(\cdot, t)\|_{L^2(0,1)}^2} - \underbrace{u_x(x, t) u(x, t)}_{\substack{0 \text{ by BC}}} \Big|_{x=0}^1 + \underbrace{\int_0^1 u_x(x, t) u_x(x, t) dx}_{\|u_x(\cdot, t)\|_{L^2(0,1)}^2} =$$

Hence, we get:

$$\frac{1}{2} \frac{d}{dt} \|u(\cdot, t)\|_{L^2(0,1)}^2 + \underbrace{\|u_x(\cdot, t)\|_{L^2(0,1)}^2}_{\geq 0} = \underbrace{\int_0^1 f(x, t) u(x, t) dx}_{C-S} \quad (*)$$

Using C-S, we get:

$$\underbrace{\frac{1}{2} \frac{d}{dt} \|u(\cdot, t)\|_{L^2(0,1)}^2}_{\text{Chain rule}} \leq \|f(\cdot, t)\|_{L^2(0,1)} \cdot \|u(\cdot, t)\|_{L^2(0,1)}$$

$$\frac{2}{2} \|u(\cdot, t)\|_{L^2} \cdot \frac{d}{dt} \|u(\cdot, t)\|_{L^2(0,1)}$$

This gives

$$\frac{d}{dt} \|u(\cdot, t)\|_{L^2(0,1)} \leq \|f(\cdot, t)\|_{L^2(0,1)}$$

Finally, integrate in time: $\int_0^t ds$

$$\|u(\cdot, t)\|_{L^2(0,1)} - \underbrace{\|u(\cdot, 0)\|_{L^2}}_{\|u_0\|_{L^2(0,1)}} \leq \int_0^t \|f(\cdot, s)\|_{L^2(0,1)} ds$$

$\therefore$

(ii) $\leadsto$ book if interested

(iii) We assume $f \equiv 0$ and consider (*):

$$\frac{1}{2} \frac{d}{dt} \|u(\cdot, t)\|_{L^2(\Omega)}^2 + \|u_x(\cdot, t)\|_{L^2(\Omega)}^2 \leq 0$$

Remembering Poincaré (1D) $\|u(\cdot, t)\|_{L^2(\Omega)}^2 \leq \frac{1}{2} \|u_x(\cdot, t)\|_{L^2(\Omega)}^2$

We get

[need $v(0) = 0$ and not $v \in H_0^1$]

$$\frac{1}{2} \frac{d}{dt} \|u(\cdot, t)\|_{L^2(\Omega)}^2 + 2 \|u(\cdot, t)\|_{L^2(\Omega)}^2 \leq 0$$

Integrating factor $\cdot 2e^{4t}$ to get:

$$\frac{d}{dt} \left(\|u(\cdot, t)\|_{L^2(\Omega)}^2 \right) e^{4t} + 4e^{4t} \|u(\cdot, t)\|_{L^2(\Omega)}^2 \leq 0$$

$$\frac{d}{dt} \left(e^{4t} \|u(\cdot, t)\|_{L^2(\Omega)}^2 \right) \underset{\substack{\text{prod} \\ \text{rule} \\ + \text{chain}}}{=} 4e^{4t} \cdot \|u\|^2 + e^{4t} \frac{d}{dt} \|u\|^2$$

Finally, we integrate in time $\int_0^t \dots ds$:

$$e^{4t} \|u(\cdot, t)\|_{L^2(\Omega)}^2 - \|u_0\|_{L^2(\Omega)}^2 \leq 0$$

$$\hookrightarrow \|u(\cdot, t)\|_{L^2(\Omega)} \leq e^{-2t} \|u_0\|_{L^2(\Omega)} \quad \therefore)$$



Th: (Energy estimate)

Consider (H) with $f \equiv 0$ and let $\varepsilon > 0$. Then,

$$\int_{\boxed{\varepsilon}}^{\varepsilon} \|u_t(\cdot, s)\|_{L^2(0,1)}^2 ds \leq \frac{1}{2} \sqrt{\ln\left(\frac{\varepsilon}{\boxed{\varepsilon}}\right)} \|u_0\|_{L^2(0,1)}^2$$

log

for all $t \in [\varepsilon, T]$.

Proof:

• Multiply DE $-t u_{xx}$ and $\int_0^1 \dots dx$ to get:

$$-t \underbrace{\int_0^1 u_t u_{xx} dx}_{\text{by part}} + t \underbrace{\int_0^1 u_{xx} \cdot u_{xx} dx}_{= \|u_{xx}\|_{L^2}^2} = 0$$

$$\left. u_t u_x \right|_{x=0}^1 - \int_0^1 \underbrace{u_{tx} u_x}_{\frac{1}{2} \frac{d}{dt} u_x^2} dx$$

0 by BC ($u(0,t)=0, u_x(1,t)=0$)
 $u_t(0,t)=0$

We obtain:

$$\frac{t}{2} \frac{d}{dt} \|u_x\|_{L^2}^2 + t \|u_{xx}\|_{L^2}^2 = 0$$

(+)

• Next, observe that:

$$t \frac{d}{dt} \|u_x\|_L^2 = \frac{d}{dt} (t \|u_x\|_L^2) - \|u_x\|_L^2$$

product rule, $t \frac{d}{dt} \|u_x\|_L^2 + 1 \cdot \|u_x\|_L^2$

and thus (1) reads

$$\frac{1}{2} \frac{d}{dt} (t \|u_x\|_L^2) - \frac{1}{2} \|u_x\|_L^2 + t \|u_{xx}\|_L^2 = 0$$

$$\Leftrightarrow \frac{d}{dt} (t \|u_x\|_L^2) + 2t \|u_{xx}\|_L^2 = \|u_x\|_L^2$$

• Integrate in time $\int_0^t \dots ds$ and get:

$$t \|u_x\|_L^2 - 0 + 2 \int_0^t s \|u_{xx}\|_L^2 ds = \int_0^t \|u_x\|_L^2 ds \leq \frac{1}{2} \|u_0\|_L^2 \quad (*)$$

[Show (*):

From previous proof: $\frac{1}{2} \frac{d}{ds} \|u(\cdot, s)\|_L^2 + \|u_x(\cdot, s)\|_L^2 = 0$

integrate

$$\int_0^t \left(\frac{1}{2} \frac{d}{ds} \|u(\cdot, s)\|_L^2 + \|u_x(\cdot, s)\|_L^2 \right) ds = 0$$

$$\Leftrightarrow \frac{1}{2} \|u(\cdot, t)\|_L^2 - \frac{1}{2} \|u_0\|_L^2 + \int_0^t \|u_x(\cdot, s)\|_L^2 ds = 0$$

$$\Rightarrow \int_0^t \|u_x(\cdot, s)\|_{L^2}^2 ds = \frac{1}{2} \|u_0\|_{L^2}^2 - \underbrace{\frac{1}{2} \|u(\cdot, t)\|_{L^2}^2}_{\geq 0} \leq \frac{1}{2} \|u_0\|_{L^2}^2$$

This (*)

The above provides:

$$t \|u_x(\cdot, t)\|_{L^2}^2 + 2 \int_0^t s \|u_{xx}(\cdot, s)\|_{L^2}^2 ds \leq \frac{1}{2} \|u_0\|_{L^2}^2$$

which implies

$$2 \int_0^t s \|u_{xx}(\cdot, s)\|_{L^2}^2 ds \leq \frac{1}{2} \|u_0\|_{L^2}^2$$

$$\Rightarrow \left(\int_0^t s \|u_{xx}(\cdot, s)\|_{L^2}^2 ds \right)^{1/2} \leq \frac{1}{2} \|u_0\|_{L^2} \quad (\text{X})$$

Finally, we use again DE $u_t - u_{xx} = 0 \Rightarrow$

$u_t = u_{xx} \Rightarrow \|u_t\|_{L^2(0,1)} = \|u_{xx}\|_{L^2(0,1)}$ and integrate in time:

$$\int_{\varepsilon}^t \|u_t(\cdot, s)\|_{L^2(0,1)}^2 ds = \int_{\varepsilon}^t \|u_{xx}(\cdot, s)\|_{L^2(0,1)}^2 ds \stackrel{C-5}{\leq} \int_{\varepsilon}^t \frac{1}{\sqrt{s}} ds$$

$$\leq \left(\int_{\varepsilon}^t \frac{1}{(\sqrt{s})^2} ds \right)^{1/2} \cdot \left(\int_{\varepsilon}^t \|u_{xx}(\cdot, s)\|_{L^2}^2 \cdot s ds \right)^{1/2}$$

exact

$$\lesssim \frac{1}{2} \sqrt{\ln\left(\frac{t}{\varepsilon}\right)} \cdot \underbrace{\|u_0\|_2}_{\sim}$$

