

$\tau \rightarrow$ source / sink
 (heat / window)
 $\rightarrow$ convection / advection
 (transport of a chemical in fluid)

(13)

Let us finally consider another type of BC:

Solve the following BVP with $C(1)$ FEM;

$$(BVP) \begin{cases} -a u''(x) + b u'(x) = \tau & \text{for } 0 < x < 1 \\ u(0) = 0, u'(1) = \beta \neq 0 \end{cases}$$

Here $a, b > 0, \tau, \beta \neq 0$ are given, wlog $a = b = \tau = 1$

Rem. Homogeneous Dirichlet BC in $x=0$

(P)

Inhomogeneous Neumann BC in $x=1$ ($u'(1) = \beta \neq 0$)

Prescribe FLUX at boundary

(i) To get a VF, consider the space $V = \{u: (0,1) \rightarrow \mathbb{R} : u, u' \in L^2(0,1), u(0) = 0\}$
Dirichlet

Then, multiply DE by a test fct $v \in V$ and integrate:

$$-a \int_0^1 u'(x) v'(x) dx + b \int_0^1 u'(x) v(x) dx = \tau \int_0^1 v(x) dx \quad (\Rightarrow)$$

$$-a u'(x) v(x) \Big|_0^1 + a \int_0^1 u'(x) v'(x) dx + b \int_0^1 u'(x) v(x) dx = \tau \int_0^1 v(x) dx$$

$$-a u'(1) v(1) + a u'(0) v(0) \quad \begin{matrix} \beta \text{ " since} \\ \text{Neumann} \end{matrix} \quad \begin{matrix} "0 \text{ since } v \in V \end{matrix}$$

We get:

$$(VF) \text{ Find } u \in V \text{ s.t. } a(u', v')_{L^2} + b(u', v)_{L^2} = \tau \int_0^1 v(x) dx + a\beta u(1) \quad \forall v \in V.$$

(ii) To get a FE problem consider

Uniform partition $\mathcal{T}_h: 0 = x_0 < x_1 < \dots < x_{m+1} = 1$ with $h = x_{j+1} - x_j = \frac{1}{m+1}$

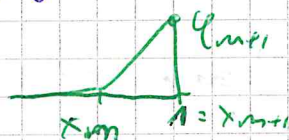
Space $V_h = \{u \in C(0,1) \rightarrow \mathbb{R} : u \text{ is cont. plus linear on } T_h \text{ and } u(0) = 0\} = \text{span}(\varphi_1, \varphi_2, \dots, \varphi_{m+1})$

We get

$$(FE) \text{ Find } u_h \in V_h \text{ s.t. } a(u_h', X')_{L^2} + b(u_h', X)_{L^2} = \tau \int_0^1 X(x) dx + a\beta X(1) \quad \forall X \in V_h$$

(iii) To get the linear syst. of eq., choose $X = \varphi_i, i=1,2,\dots,m+1$ and write $u_h(x) = \sum_{j=1}^{m+1} \zeta_j \varphi_j(x)$ in (FE).

! φ_{m+1} is half hat



The FE problem then reads:

$$a \cdot \left(\sum_{j=1}^{m+1} \zeta_j \varphi_j', \varphi_i' \right)_{L^2} + b \cdot \left(\sum_{j=1}^{m+1} \zeta_j \varphi_j', \varphi_i \right)_{L^2} = \tau \int_0^1 \varphi_i(x) dx + a\beta \varphi_i(1) \quad \text{for } i=1,2,\dots,m+1$$

($\Rightarrow$) (linearity)

$$a \sum_{j=1}^{m+1} \underbrace{(\varphi_j', \varphi_i')}_S \cdot \tilde{z}_j + b \sum_{j=1}^{m+1} \underbrace{(\varphi_j', \varphi_i')}_C \cdot \tilde{z}_j = \underbrace{2 \int_0^1 \varphi_i(x) dx}_R + a \beta \varphi_i(1)$$

$$\Leftrightarrow a \cdot S \tilde{z} + b C \tilde{z} = R \Leftrightarrow (aS' + bC) \tilde{z} = R \quad (\text{lin. syst. of eq.})$$

(5) $S \Leftarrow$ stiffness matrix
($m+1$) $\times$ ($m+1$)

$$S = \frac{1}{h} \begin{pmatrix} 2 & -1 & & 0 \\ -1 & 2 & -1 & \\ & -1 & 2 & -1 \\ 0 & & -1 & 2 \end{pmatrix}$$

$\Delta \varphi_{m+1}$ half hat

$C \Leftarrow$ convection matrix
($m+1$) $\times$ ($m+1$)

$$C = \frac{1}{2} \begin{pmatrix} 0 & 1 & & 0 \\ -1 & 0 & 1 & \\ & 0 & -1 & 0 \\ 0 & & 0 & 1 \end{pmatrix}$$

$$R = \begin{pmatrix} \pi h \\ \pi h \\ \vdots \\ \pi h \\ \frac{\pi h}{2} + a\beta \end{pmatrix} \quad (m+1) \times 1$$

half hat

Neumann BC

Rem: Stiffness and mass matrices are symmetric $\therefore$
Convection matrix is not $\therefore$

3) Equivalence between (BVP) and (VF):

We consider the problem (if cont. + bounded)

$$(BVP) \quad \begin{cases} -u''(x) = f(x) & \text{for } 0 < x < 1 \\ u(0) = u(1) = 0 \end{cases} \quad \begin{matrix} (BC) \\ (BC) \end{matrix}$$

and show that it is equivalent to

$$(VF) \text{ Find } u \in V^0 \text{ s.t. } (u', v')_{L^2(0,1)} = (f, v)_{L^2(0,1)} \quad \forall v \in V^0$$

if u is twice cont. diff. ($u \in C^2(0,1)$)

$$\text{Recall: } V^0 = H_0^1(0,1) = \{v \in C^1(0,1) : v, v' \in L^2(0,1) \text{ and } v(0) = v(1) = 0\}$$

Th: (BVP) $\Leftrightarrow$ (VF) ~~then~~ and $u \in C^2(0,1)$

Proof:

$\Rightarrow$ Already done; multiply with test φ and integrate

$$\int_0^1 u''(x) v(x) dx = \int_0^1 f(x) v(x) dx \Leftrightarrow (u', v')_0 = (f, v)_0 \quad \forall v \in V^0$$

by part

$\Leftarrow$: (VP) reads $(u', v') = (f, v) \quad \forall v \in U$ on either side by parts

$$-\int_0^1 u''(x) v(x) dx + \underbrace{u'(x) v(x)}_{=0 \text{ since } v \in U^0} \Big|_0^1 = \int_0^1 f(x) v(x) dx \quad \text{on}$$

$$\int_0^1 (u''(x) + f(x)) v(x) dx = 0 \quad \forall v \in U^0 \quad (*)$$

We show by contradiction that this implies (BVP) $u''(x) + f(x) = 0$ on $(0, 1)$.

Suppose $u'' + f \neq 0$. Then, by continuity, $\exists x \in (0, 1)$ s.t.

$u'' + f \neq 0$ in a neighborhood of x . If we choose now $v \in U^0$ that has the same sign in this neighborhood and is zero everywhere else, we obtain $\int_0^1 (u''(x) + f(x)) v(x) dx > 0$.

This is a contradiction to $(*)$!!

Therefore, we must have $u''(x) + f(x) = 0$ on $(0, 1)$, i.e. (BVP) $\square$

Def. We denote by $H_0^1(a, b)$ the space

$$H_0^1(a, b) = \{v: [a, b] \rightarrow \mathbb{R}; v, v' \in L^2(a, b), v(a) = v(b) = 0\}$$

4) Error estimates:

We start with

⑦

th. (Poincaré inequality)

Let $L > 0$ and the open interval $\mathcal{R} = (0, L)$. Assume $u \in H_0^1(\mathcal{R})$, then

$$\|u\|_{L^2(\mathcal{R})} \leq C_L \cdot \|u'\|_{L^2(\mathcal{R})}$$

Proof:

Since $u \in H_0^1(\mathcal{R})$, then $u(0) = 0$ and $u(x) = \int_0^x u'(y) dy \quad \forall x \in \mathcal{R}$.

It follows:

$$|u(x)|^2 = \left| \int_0^x 1 \cdot u'(y) dy \right|^2 \stackrel{\text{CS}}{\leq} \left| \int_0^x 1 dy \right| \cdot \int_0^x |u'(y)|^2 dy \stackrel{x \leq L}{\leq} x \cdot \int_0^L |u'(y)|^2 dy$$

Finally, integrate w.r.t x and get

$$\underbrace{\int_0^L |u(x)|^2 dx}_{\|u\|_{L^2(0,1)}^2} \leq \underbrace{\int_0^L x dx}_{\frac{L^2}{2}} \cdot \int_0^L |u'(y)|^2 dy \quad \text{on} \quad \|u\|_{L^2(0,1)}^2 \leq \frac{L^2}{2} \cdot \|u'\|_{L^2(0,1)}^2$$

$\stackrel{C_L^2 \text{ (indep. of } u)}}{\|u\|_{L^2(0,1)}^2} \quad \square$

⑤

Can bound u only with L and u'

Def. For $f, g \in H_0^1(a, b)$, we define the energy inner product

$$(f, g)_E = \int_a^b f'(x)g'(x) dx = (f', g')_{L^2(a, b)}$$

and the energy norm

$$\|f\|_E = \sqrt{(f, f)_E} = \sqrt{(f', f')_{L^2(a, b)}} = \|f'\|_{L^2(a, b)}$$

The above definition is used to study error of FEM from

$$(BVP) \quad \begin{cases} -u''(x) = f(x) & \text{for } x \in (0, 1) \\ u(0) = u(1) = 0 \end{cases}$$

Recall,

(VF) Find $u \in H_0^1(0, 1)$ s.t. $(u', v')_L = (f, v)_L \quad \forall v \in H_0^1(0, 1)$

(FE) Find $u_h \in V_h^0$ s.t. $(u_h', X')_L = (f, X)_L \quad \forall X \in V_h^0$, where

Uniform partition $T_h: 0 = x_0 < x_1 < \dots < x_m < x_{m+1} = 1, x_{j+1} - x_j = h$

$V_h^0 = \text{span}(\psi_1, \psi_2, \dots, \psi_m)$, ψ_j hat fun.

Obs $V_h^0 \subset H_0^1$ so that one can take $v \in X$ in (VF)! And get

$$(u', X') = (f, X')$$

$$(u_h', X') = (f, X')$$

$$\Rightarrow \underline{(u' - u_h', X') = 0 \quad \forall X \in V_h^0!}$$

This is known as Galerkin orthogonality

Rem: The above means that the error $e = u - u_h$ of FEM is orthogonal to V_h^0 in the energy inner product.

This is now used to show

Th: Let u be the sol. to (BVP) and u_h the sol. to (FE),

One has

$$\|u - u_h\|_E \leq \|u - X\|_E \quad \forall X \in V_h^0$$

That is the FE sol. is the best approx of u in V_h^0 in energy norm.

Proof:

Consider Def norm

$X \in V_h^0$

Linearity

$$\begin{aligned} \|u - u_h\|_E^2 &= (u' - u_h', u' - u_h')_L \stackrel{\text{Def norm}}{=} (u' - u_h', u' - X' + X' - u_h')_L \stackrel{\text{Linearity}}{=} \\ &= (u' - u_h', u' - X') + (u' - u_h', X' - u_h') \stackrel{\text{Galerkin orthog.}}{=} \|u' - u_h'\|_E \cdot \|u' - X'\|_E \quad (\because) \end{aligned}$$