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MVE555 Architectural Geometry, Lecture 2



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Mathematical Sciences

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19th September 2019



Linear Transformations

Rigid Body Motion

Homogeneous Coordinates

Projections

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Linear Transformations — Matrices

Linear transformations on \mathbb{R}^3 are studied in linear algebra, and are characterised by linearity:

$$\begin{cases} T(\boldsymbol{x} + \boldsymbol{y}) = T(\boldsymbol{x}) + T(\boldsymbol{y}), & \forall \boldsymbol{x}, \boldsymbol{y} \in \mathbb{R}^3 \\ T(\lambda \boldsymbol{x}) = \lambda T(\boldsymbol{x}), & \forall \lambda \in \mathbb{R}, \boldsymbol{x} \in \mathbb{R}^3. \end{cases}$$

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If $\boldsymbol{x} = (x_1, x_2, x_3)$ then $\boldsymbol{y} = (y_1, y_2, y_3) = T(\boldsymbol{x})$ can be written as a matrix multiplication:

$$\begin{bmatrix} y_1 \\ y_2 \\ y_3 \end{bmatrix} = \underbrace{\begin{bmatrix} a_{11} & a_{12} & a_{13} \\ a_{21} & a_{22} & a_{23} \\ a_{31} & a_{32} & a_{33} \end{bmatrix}}_{\boldsymbol{A}} \begin{bmatrix} x_1 \\ x_2 \\ x_3 \end{bmatrix},$$

where A is constant.

Images of the Basic Unit Vectors

Since (1,0,0) is transformed into

$$\begin{bmatrix} a_{11} & a_{12} & a_{13} \\ a_{21} & a_{22} & a_{23} \\ a_{31} & a_{32} & a_{33} \end{bmatrix} \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} a_{11} \\ a_{21} \\ a_{31} \end{bmatrix},$$

the first column of A is the destination of (1, 0, 0).

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the first column of ${\bm A}$ is the destination of (1,0,0). Similarly, the second and third column tell us where (0,1,0) and (0,0,1) go!



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- If A is an orthogonal matrix, i.e. $A^{\intercal}A = I$, then $A^{-1} = A^{\intercal}$
- Whatever $oldsymbol{A}$ is, the origin is never moved, i.e. $oldsymbol{A0}=oldsymbol{0}$

Definition of Rigid Body Motion

A rigid body motion is composed of a *rotation* and a *translation*:



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If combined with a scaling, it becomes a similarity transformation.

Let R be a rotation matrix (later slides) and t be a vector. Then

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- Rigid body motions are associative
- Not a linear transformation the origin is moved!
- We will see later how to write them using only a matrix multiplication anyway!

Representing Rotations — 2D

In 2D, rotations are almost always represented using the 2×2 matrix

$$\boldsymbol{R}(\varphi) = \begin{bmatrix} \cos\varphi & -\sin\varphi\\ \sin\varphi & \cos\varphi \end{bmatrix},$$

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To rotate a point ${\boldsymbol x} = (x,y)$ and angle φ about the origin, we do

$$\boldsymbol{y} = \boldsymbol{R}(\varphi)\boldsymbol{x} = \begin{bmatrix} \cos\varphi & -\sin\varphi \\ \sin\varphi & \cos\varphi \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} x\cos\varphi - y\sin\varphi \\ x\sin\varphi + y\cos\varphi \end{bmatrix}$$

In 3D, rotations around the three coordinate axes are written as

$$\boldsymbol{R}_{x}(\alpha) = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos \alpha & -\sin \alpha \\ 0 & \sin \alpha & \cos \alpha \end{bmatrix},$$
$$\boldsymbol{R}_{y}(\beta) = \begin{bmatrix} \cos \beta & 0 & \sin \beta \\ 0 & 1 & 0 \\ -\sin \beta & 0 & \cos \beta \end{bmatrix},$$
$$\boldsymbol{R}_{z}(\gamma) = \begin{bmatrix} \cos \gamma & -\sin \gamma & 0 \\ \sin \gamma & \cos \gamma & 0 \\ 0 & 0 & 1 \end{bmatrix}.$$

Representing Rotations — 3D

• Any 3D rotation can be obtained using Tait-Bryan angles as

 $\boldsymbol{R}(\alpha,\beta,\gamma) = \boldsymbol{R}_x(\alpha)\boldsymbol{R}_y(\beta)\boldsymbol{R}_z(\gamma)$

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- By performing different combinations of the rotations, we get various *Euler angle* representations no clear standard!
- It is often easier to think using an *axis-angle representation*, e.g. *Rodrigues' formula*:

$$\boldsymbol{R} = \boldsymbol{I} + \sin \varphi [\boldsymbol{v}]_{\times} + (1 - \cos \varphi) [\boldsymbol{v}]_{\times}^2$$

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Rodrigues' Formula



Rodrigues' Formula Proof

We have $x = x_{\parallel} + x_{\perp}$, where x_{\parallel} is parallel to v (and thus does not change), and x_{\perp} is perpendicular to v. Note also that x_{\perp} and $v \times x_{\perp}$ make up an orthogonal basis in the plane orthogonal to v. It follows that

$$\begin{split} \boldsymbol{R} \boldsymbol{x}_{\perp} &= \cos \varphi \; \boldsymbol{x}_{\perp} + \sin \varphi \; (\boldsymbol{v} \times \boldsymbol{x}_{\perp}) \\ &= -\cos \varphi \; (\boldsymbol{v} \times (\boldsymbol{v} \times \boldsymbol{x}_{\perp})) + \sin \varphi \; (\boldsymbol{v} \times \boldsymbol{x}_{\perp}) \\ &= -\cos \varphi \; (\boldsymbol{v} \times (\boldsymbol{v} \times \boldsymbol{x})) + \sin \varphi \; (\boldsymbol{v} \times \boldsymbol{x}). \end{split}$$

Thus

$$\begin{aligned} \boldsymbol{R} \boldsymbol{x} &= \boldsymbol{R} \boldsymbol{x}_{\parallel} + \boldsymbol{R} \boldsymbol{x}_{\perp} \\ &= (\boldsymbol{v}^{\intercal} \boldsymbol{x}) \boldsymbol{v} - \cos \varphi \left(\boldsymbol{v} \times (\boldsymbol{v} \times \boldsymbol{x}) \right) + \sin \varphi \left(\boldsymbol{v} \times \boldsymbol{x} \right) \\ &= \boldsymbol{x} + \sin \varphi \left(\boldsymbol{v} \times \boldsymbol{x} \right) + (1 - \cos \varphi) (\boldsymbol{v} \times (\boldsymbol{v} \times \boldsymbol{x})). \end{aligned}$$

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The Planar Case

Suppose we are working in the plane, and have a point (x, y). The plane can be 'embedded' in 3D as the plane z = 1:



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The Planar Case (contd.)

• Each point in the plane $z=1\ {\rm corresponds}$ to a 3D-line through the origin

Möbius, Der barycentrische Calcul - ein neues Hülfsmittel zur analytischen Behandlung der Geometrie, 1827.

- Each point in the plane z = 1 corresponds to a 3D-line through the origin
- The line through (x,y,1) includes $(\lambda x,\lambda y,\lambda)$ for any scalar λ

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- When $\lambda \to \pm \infty$ we obtain *ideal points*, (x, y, 0), infinitely far away (on the *line at infinity*)
- This can be used to capture the difference between vectors and points!

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• Similarly to the 2D case, we add an extra coordinate that is equal to one, i.e. the homogeneous coordinates for (x, y, z) become (x, y, z, 1) (or $(\lambda x, \lambda y, \lambda z, \lambda)$ for any $\lambda \neq 0$).

The 3D Case

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- The homogeneous coordinates (x, y, z, 0) represent the point infinitely far away in the direction (x, y, z)

Revisiting Rigid Body Motions

Recall that a rigid body motion consisting of the rotation R and the translation t is written as y = Rx + t.

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Recall that a rigid body motion consisting of the rotation $m{R}$ and the translation $m{t}$ is written as $m{y}=m{R}x+m{t}$. As it turns out,

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 $\begin{bmatrix} \mathbf{J} \\ 1 \end{bmatrix} = \underbrace{\begin{bmatrix} \mathbf{0}^{\mathsf{T}} & 1 \end{bmatrix}} \begin{bmatrix} 1 \end{bmatrix}.$

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If we use homogeneous coordinates, we can represent a rigid body motion as the matrix $oldsymbol{A}$ above.

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The Pinhole Perspective Camera



A 3D point (X, Y, Z) is thus projected to (fX/Z, fY/Z, f) in the image plane — we may omit the last coordinate:

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Using homogeneous coordinates, we can write the projection as a matrix multiplication:

$$\begin{bmatrix} fX\\fY\\Z \end{bmatrix} = \begin{bmatrix} f & 0 & 0 & 0\\ 0 & f & 0 & 0\\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} X\\Y\\Z\\1 \end{bmatrix} = \underbrace{\begin{bmatrix} f & 0 & 0\\ 0 & f & 0\\ 0 & 0 & 1 \end{bmatrix}}_{K} \begin{bmatrix} I & \mathbf{0} \end{bmatrix} \begin{bmatrix} X\\Y\\Z\\1 \end{bmatrix}$$

• A camera positioned at *t* instead of the origin, and rotated a rotation *R*, is represented by the matrix

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 For us, the *focal length* f is not particularly interesting most of the time — we can set it to f = 1 for simplicity and skip K entirely

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Orthographic Cameras — Illustration



Projections — Orthographic Cameras

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