

### 1.3 Group theory 1: The classical matrix groups

Consider the **set of all** complex  $N \times N$  matrices  $U$  satisfying the unitarity condition

$$U^\dagger = U^{-1}, \text{ or } U^\dagger U = U U^\dagger = \mathbf{1}, \quad (1.5)$$

leaving the scalar product in  $\mathbf{C}^N$  invariant: If  $\chi \in \mathbf{C}^N$  and  $\chi' = U\chi$  then  $\chi'^\dagger \chi' = \chi^\dagger \chi$ . The unitarity property survives matrix multiplication which means that if two matrices  $U_1$  and  $U_2$  are unitary so is the product  $U_3 = U_1 U_2$  which hence also belongs to this matrix set:

$$U_3^\dagger U_3 = (U_1 U_2)^\dagger U_1 U_2 = U_2^\dagger U_1^\dagger U_1 U_2 = \mathbf{1}. \quad (1.6)$$

Note that the unit matrix  $\mathbf{1}$  is in this set of matrices and that every matrix has an inverse, which is actually part of the assumption above (since  $U^\dagger = U^{-1}$  implies that  $\det U = e^{i\alpha}$  with  $\alpha \in \mathbf{R}$ ). We may also add the requirement that the matrices have unit determinant which is also preserved by matrix multiplication since  $\det(U_1 U_2) = (\det U_1)(\det U_2)$ . This is a subset of the previous one without the unit determinant condition.

If we consider **all** matrices in either one of these two sets the following properties are trivially satisfied:

1. The set is **closed** under multiplication (here matrix multiplication)
2. The multiplication is **associative** (as matrix multiplication always is)
3. There is a **unit** (here the unit matrix  $\mathbf{1}$ )
4. Every element in the set has an **inverse** (true here since we consider only matrices satisfying  $U^\dagger = U^{-1}$ , i.e., matrices with non-zero determinant)

Viewing these properties instead as axioms they define a **group**, often denoted  $G$  and members of the set are called *elements*, which in the case above is called the *unitary* group:

$$U(N), \quad (1.7)$$

realised in this discussion in terms of complex  $N \times N$  matrices. When the condition that the matrices have unit determinant  $\det U = \mathbf{1}$  is added the group is called *special unitary*:

$$SU(N). \quad (1.8)$$

The standard model of elementary particles is based on three such groups:  $U(1)$ ,  $SU(2)$  and  $SU(3)$ . The group  $U(1)$  is the group of *phases*, i.e., multiplication by  $e^{i\alpha}$  where  $\alpha$  is a (real) parameter (angle), and is *abelian*. The other two groups are *non-abelian*, i.e.,  $g_1 g_2 \neq g_2 g_1$  for some  $g_1, g_2 \in G$ . Any  $U \in SU(2)$  can be written in terms of  $a, b \in \mathbf{C}$  as

$$U \in SU(2) : U = \begin{pmatrix} a & b \\ -b^* & a^* \end{pmatrix}, \text{ where } |a|^2 + |b|^2 = 1, \quad (1.9)$$

which is easily checked. Two general such matrices do not commute and the group is therefore *non-abelian*. The set of matrices in this group can be parametrised by the points

on the unit three-sphere  $S^3$  since  $|a|^2 + |b|^2 = 1$  which, if written in real variables  $a = x_1 + ix_2$  and  $b = x_3 + ix_4$ , becomes  $x_1^2 + x_2^2 + x_3^2 + x_4^2 = 1$ , i.e., the equation for  $S^3$  embedded in  $\mathbf{R}^4$ .

While  $U(1)$  and  $SU(2)$  can be viewed in this way in terms of simple geometries  $S^1$  (unit complex numbers) and  $S^3$  (unit quaternions)<sup>3</sup>  $SU(3)$  and higher  $SU(N)$  groups are much more complicated as manifolds and do not correspond to anything with other names.

These notions can be carried over to real matrices leading to the *orthogonal* groups

$$O(N) \text{ or } SO(N). \quad (1.10)$$

Among the groups called *classical* there is just one other case, the *symplectic* groups. One way to define matrix groups is by looking for matrices preserving some special matrices numerically. One example is the unit matrix which then leads to orthogonal groups ( $g\mathbf{1}g^T = \mathbf{1} \Rightarrow g \in O(N)$ ) while invariance of the antisymmetric  $2N \times 2N$  matrix

$$C = \begin{pmatrix} 0 & \mathbf{1} \\ -\mathbf{1} & 0 \end{pmatrix}, \quad (1.11)$$

defines the *symplectic* groups<sup>4</sup>

$$Sp(2N). \quad (1.12)$$

#### Comments:

1. The complete classification of all *finite-dimensional Lie groups* (i.e., continuous of the kind discussed above) is known as the Cartan classification and contains in addition to the *classical* groups above, i.e., the unitary (denoted  $A_n$  by Cartan), orthogonal (denoted  $B_n$  or  $D_n$ ), and symplectic (denoted  $C_n$ ), also the *exceptional* ones  $G_2, F_4, E_6, E_7, E_8$ . These latter ones have, however, no simple definition in terms of matrices similar to the one used above. The index  $n$  (and the indices appearing on the exceptional groups) is the **rank** of the group related to the maximal number of matrices that can be diagonalised simultaneously. The classes  $A_n$  contain also the groups  $GL(N)$  and  $SL(N)$  which are general matrices with non-zero or unit determinant, respectively.
2. There are also other important groups like *finite* ones with a finite number of group elements, and those with a *discrete* set of elements which is *infinite* in number (see courses in Group theory and in String theory).
3. Other important Lie groups in physics have *infinite dimension*. Examples of such are *Virasoro* and *Kac-Moody* appearing in two-dimensional *conformal field theory* ( $CFT_2$ ) used, e.g., in string theory and in the context of phase transitions in condensed matter systems.
4. There is a very important (also for QFT) distinction between compact groups (e.g.,  $U(N), SU(N), SO(N)$ ) and non-compact ones (e.g.,  $SO(1,3), SU(1,1), Sp(2N), SL(N)$ ) discussed further in "Group theory 2" on Lie algebras and representations.

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<sup>3</sup>There is a third case, the octonions, for which a unit octonion can be shown to be the same as the seven-dimensional sphere,  $S^7$ . This manifold is, however, not a group manifold but does nevertheless share some properties with group manifolds as, e.g., being parallelisable.  $S^7$  plays a key role in string/M theory.

<sup>4</sup>Note that some authors call these groups  $Sp(N)$ .